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## On the Essential Self-Adjointness of Anti-Commutative Operators

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**Abstract.** In this article, linear operators satisfying anti-commutation relations are considered. It is proven that an anti-commutative type of the Glimm-Jaffe-Nelson commutator theorem follows.

## 1 Introduction and Main Theorem

In this article we consider the self-adjointness of linear operators satisfying anti-commutation relations. For criteria on the self-adjointness of the symmetric operator satisfying a commutation relation, there is the Glimm-Jaffe-Nelson commutator theorem. Refer to [2, 3, 4] of the original papers, and see also e.g. ([1]; Theorem 2.32, [6]; Theorem X.36). We investigate an anti-commutative type of the commutator theorem.

### Theorem 1

*Let  $H$  be a symmetric operator and  $A$  be a strictly positive self-adjoint operator i.e. there exists a constant  $\delta_A > 0$  such that  $(\Psi, A\Psi) \geq \delta_A(\Psi, \Psi)$ . We assume the following conditions :*

*(C.1) There exists a core  $\mathcal{D}_0$  of  $A$  such that  $\mathcal{D}_0 \subset \mathcal{D}(H)$  where  $\mathcal{D}(H)$  denotes the domain of  $H$ .*

*(C.2) There exists a constant  $a > 0$  such that for all  $\Psi \in \mathcal{D}_0$ ,*

$$\|H\Psi\| \leq a\|A\Psi\|. \quad (1)$$

*(C.3) There exists a constant  $b > 0$  such that for all  $\Psi \in \mathcal{D}_0$ ,*

$$\left| (H\Psi, A\Psi) + (A\Psi, H\Psi) \right| \leq b\|A^{1/2}\Psi\|^2. \quad (2)$$

*Then  $H$  is essentially self-adjoint on  $\mathcal{D}_0$ .*

**Remark 1** Let  $X$  and  $Y$  be symmetric operators on a Hilbert space  $\mathcal{X}$ . Then the real part and the imaginary part of the inner product  $(X\Psi, Y\Psi)$  for  $\Psi \in \mathcal{D}(XY) \cap \mathcal{D}(YX)$  is expressed by using the commutator and the anti-commutator :

$$\operatorname{Re}(X\Psi, Y\Psi) = \frac{(\Psi, \{X, Y\}\Psi)}{2}, \quad (3)$$

and

$$\operatorname{Im}(X\Psi, Y\Psi) = \frac{(\Psi, [X, Y]\Psi)}{2}, \quad (4)$$

where  $\{X, Y\} = XY + YX$  and  $[X, Y] = XY - YX$ . On the proof of the Glimm-Jaffe-Nelson commutator theorem, the estimate of a type of (4) is investigated. In this article we apply the estimate of a type of (3).

**Remark 2** Since  $\|A^{1/2}\Phi\| \leq \|A\Phi\| + \|\Phi\|$  follows for  $\Phi \in \mathcal{D}(A)$ , we see that  $\mathcal{D}_0$  is a core of  $A^{1/2}$ . Then from (C.2) and (C.3), it is seen that  $\mathcal{D}(A) \subset \mathcal{D}(\overline{H_{\uparrow\mathcal{D}_0}})$  and for  $\Phi \in \mathcal{D}(A)$ ,

$$\left| (\overline{H_{\uparrow\mathcal{D}_0}}\Phi, A\Phi) + (A\Phi, \overline{H_{\uparrow\mathcal{D}_0}}\Phi) \right| \leq b\|A^{1/2}\Phi\|^2 \quad (5)$$

hold, where  $\overline{X}$  denotes the closure of a operator  $X$ .

**(Proof of Theorem 1)**

Let  $c \in \mathbf{R}$  be a real number satisfying  $c > \frac{b}{2}$ , and let us set  $z = c + i \in \mathbf{C}$ . For a closable operator  $X$ , it is seen that  $\overline{X} = (X^*)^*$  and  $(\overline{X})^* = X^*$  follow. Then from the general theorem ([6], Theorem X.1) on the essential-self-adjointness of closed symmetric operator, it is enough to show that  $\dim \ker \left( H_{\uparrow\mathcal{D}_0}^* + z \right) = \dim \ker \left( H_{\uparrow\mathcal{D}_0}^* + z^* \right) = 0$ . Let  $\Psi \in \mathcal{D}_0$ . It is noted that  $A^{-\alpha}$ ,  $\alpha > 0$  is bounded, since  $A$  is a strictly positive. Then by using  $\operatorname{Re}(f, g) = \frac{(f, g) + (g, f)}{2}$  and  $(H_{\uparrow\mathcal{D}_0}^*)^* = \overline{H_{\uparrow\mathcal{D}_0}}$ , we see that for  $\Xi = A^{-1}\Psi \in \mathcal{D}(A)$ ,

$$\operatorname{Re} \left( \Xi, (H_{\uparrow\mathcal{D}_0}^* + z)\Psi \right) = \frac{1}{2} \left( (\overline{H_{\uparrow\mathcal{D}_0}}\Xi, A\Xi) + (A\Xi, \overline{H_{\uparrow\mathcal{D}_0}}\Xi) \right) + c\|A^{1/2}\Xi\|^2.$$

Then by using (5) and  $\|A^{1/2}\Xi\| = \|A^{-1/2}\Psi\|$ , we have

$$\left| \operatorname{Re} \left( \Xi, (H_{\uparrow\mathcal{D}_0}^* + z)\Psi \right) \right| \geq \left( c - \frac{b}{2} \right) \|A^{-1/2}\Psi\|. \quad (6)$$

Now let us assume that  $\Psi \in \ker \left( H_{\uparrow\mathcal{D}_0}^* + z \right)$ . Then, we see from (6) that  $A^{-1/2}\Psi = 0$ . Then by acting  $A^{1/2}$  the both side of this equation, we obtain  $\Psi = 0$ . Hence  $\ker \left( H_{\uparrow\mathcal{D}_0}^* + z \right) = \{0\}$ . Similarly, we can prove that  $\ker \left( H_{\uparrow\mathcal{D}_0}^* + z^* \right) = \{0\}$ . Thus the proof is obtained. ■

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